

Recent results on multiterminal Josephson junctions

Régis Mélin

Institut NÉEL

Centre National de la Recherche Scientifique

Université Grenoble-Alpes

France



List of Collaborators: theory

Institut Néel

Régis Mélin, Denis Feinberg

Laboratoire de Physique Théorique et des Hautes Energies
(Jussieu, Paris, France)

Benoît Douçot

Mathematics Department of INSA (Rouen, France)

Jean-Guy Caputo

List of Collaborators: experiments

Karlsruhe Institute of Technology (KIT, Germany)

Group of **Romain Danneau**

Grenoble (Néel and CEA)

Groups of **Hervé Courtois** and **François Lefloch**

Weizmann Institute (Rehovot, Israel)

Part of group of **Moty Heiblum**:

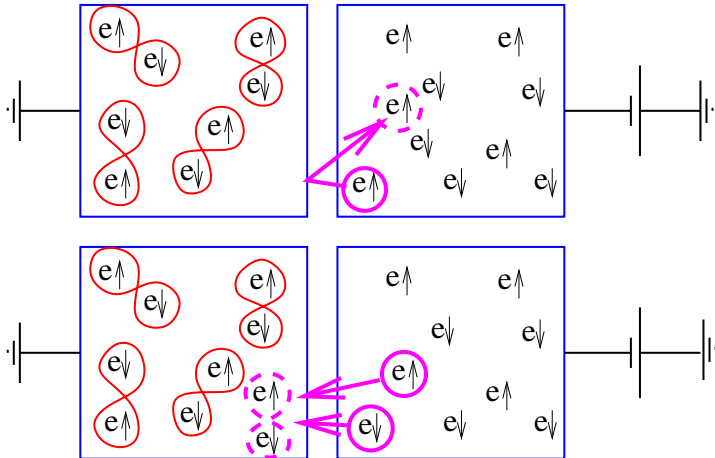
Yuval Ronen, Yonatan Cohen, Jung-Hyun Kang, Hadas Shtrikman

Harvard University

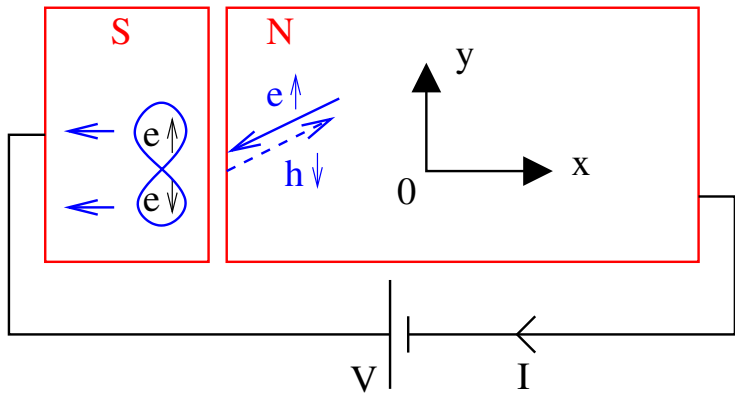
Part of group of **Philip Kim**:

Katie Huang, Yuval Ronen

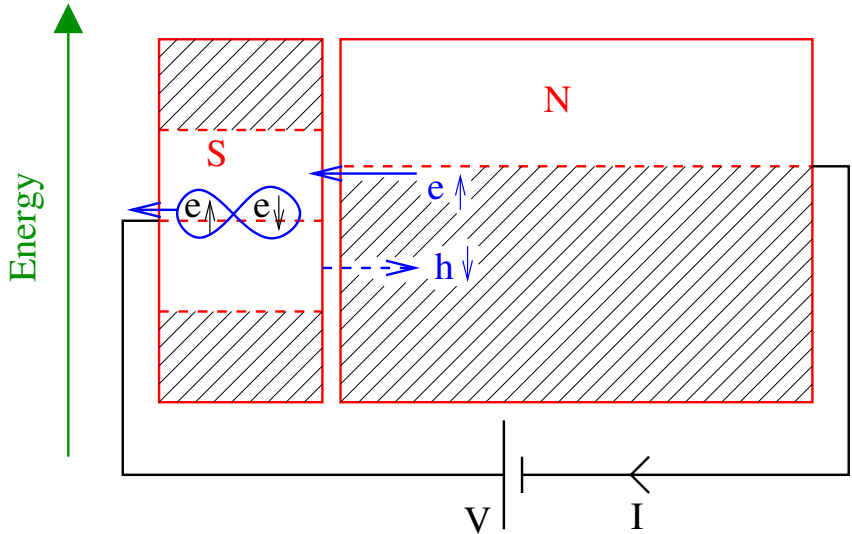
SN Junction: Andreev Reflection (1/3)



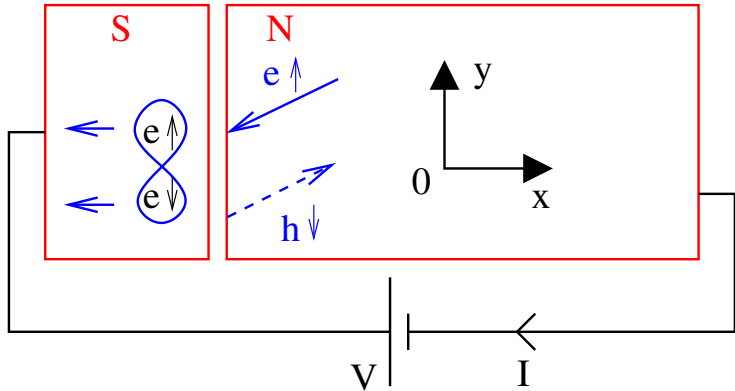
SN Junction: Andreev Reflection (2/3)



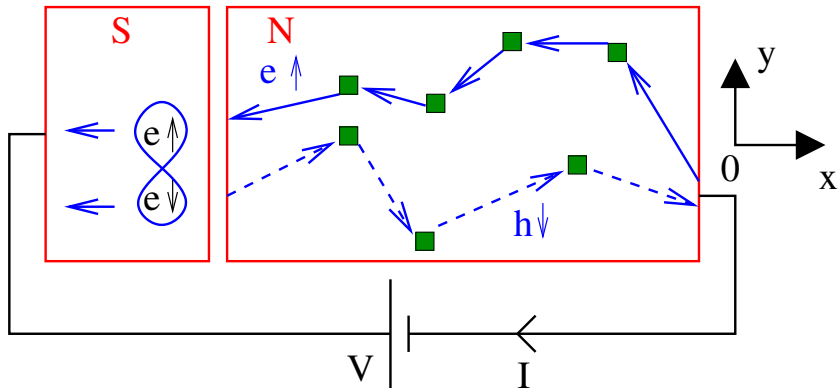
SN Junction: Andreev Reflection (3/3)



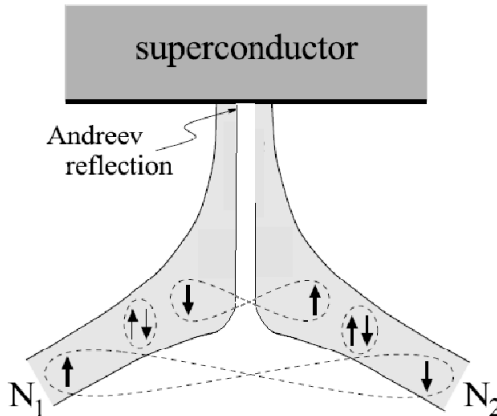
SN Junction: Nonlocal Andreev Reflection (1/3)



Nonlocal Andreev Reflection (2/3)



Nonlocal Andreev Reflection \equiv Cooper Pair Splitting (3/3)

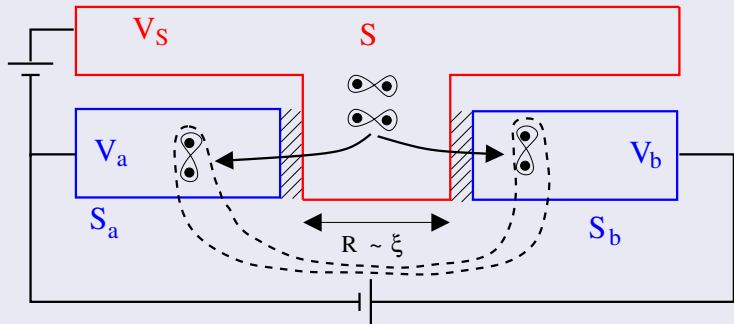


Three-terminal set-up required in experiments

First theoretical contributions: Byers-Flatté, Martin, Anatram-Datta, Deutscher-Feinberg, Falci-Hekking, Choi-Bruder-Loss, Mélin

Production of Nonlocal Quartets, $R < \xi$

A. Freyn, B. Douçot, D. Feinberg, R. Mélin, PRL 2011, NÉEL / LPTHE

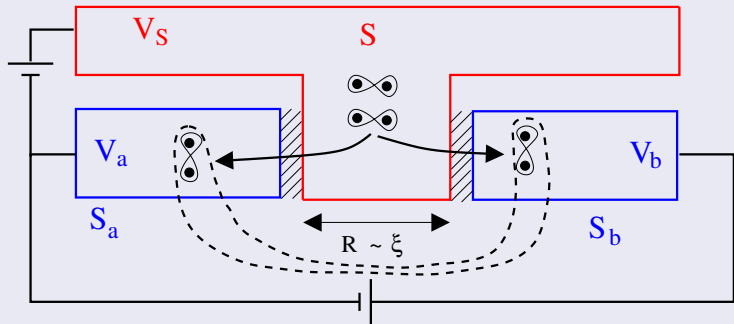


Intuitively

2 pairs in coherence volume within time interval \hbar/Δ
→ production of a correlated pair of pair between S_a and S_b

Production of Nonlocal Quartets, $R < \xi$

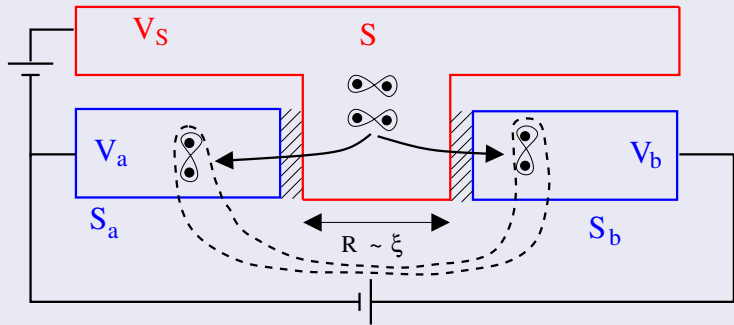
A. Freyn, B. Douçot, D. Feinberg, R. Mélin, PRL 2011, NÉEL / LPTHE



ac Josephson current of pairs from S_a to S and from S to S_b

Production of Nonlocal Quartets, $R < \xi$

A. Freyn, B. Douçot, D. Feinberg, R. Mélin, PRL 2011, NÉEL / LPTHE



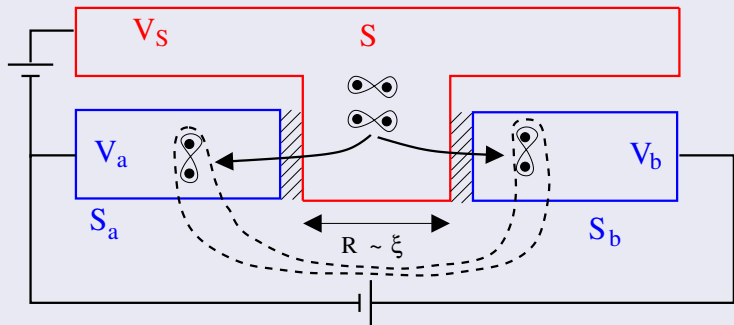
ac Josephson current of pairs from S_a to S and from S to S_b

But possibility of a dc Josephson current of quartets
from S to S_a and S_b

if $V_a = -V_b$ and $V_S = 0$ because $\Delta E = 2e(V_a + V_b - 2V_S) \equiv 0$

Production of Nonlocal Quartets, $R < \xi$

A. Freyn, B. Douçot, D. Feinberg, R. Mélin, PRL 2011, NÉEL / LPTHE



Adiabatic model

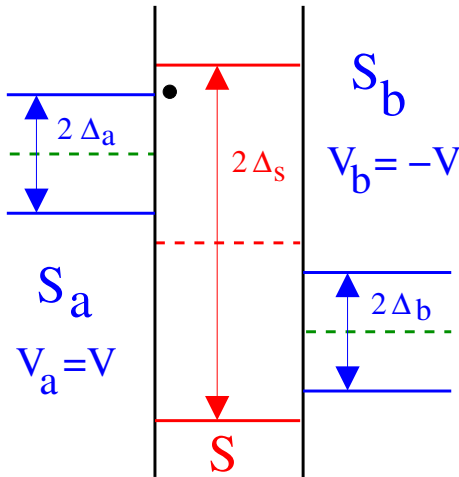
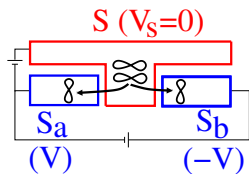
$$\phi_a(t) + \phi_b(t) - 2\phi_S = [2e(V_a + V_b - 2V_S)/\hbar]t + \phi_a + \phi_b - 2\phi_S$$
$$I_{\text{quartet}}(t) = I_c \sin[\phi_a(t) + \phi_b(t) - 2\phi_S]$$

- AC Josephson effect of quartets in general
- DC Josephson effect of quartets if $V_a = -V_b$ and $V_S = 0$

Nonlocal Quartets, DC if $V_a = -V_b$, $V_S = 0$

Quartet Resonance in DC Current for $V_a = -V_b$

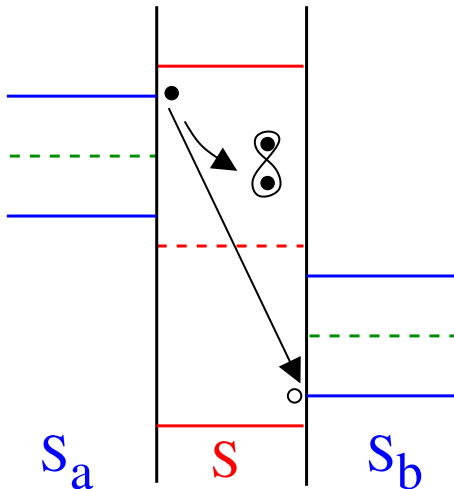
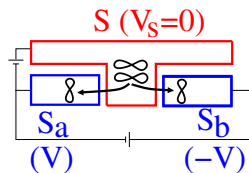
A. Freyn, B. Douçot, D. Feinberg, R. Mélin, PRL 2011, NÉEL / LPTHE



Nonlocal Quartets, DC if $V_a = -V_b$, $V_S = 0$

Quartet Resonance in DC Current for $V_a = -V_b$

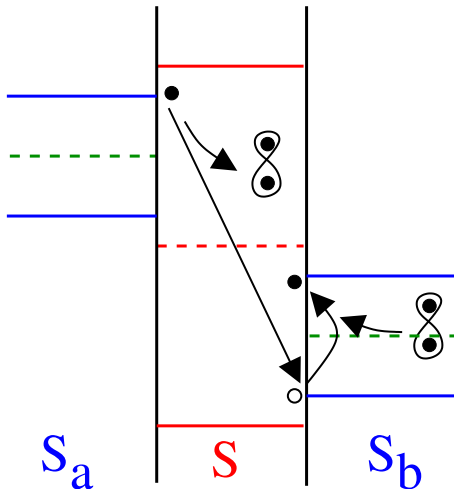
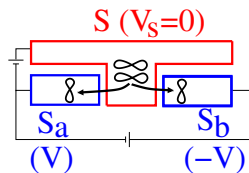
A. Freyn, B. Douçot, D. Feinberg, R. Mélin, PRL 2011, NÉEL / LPTHE



Nonlocal Quartets, DC if $V_a = -V_b$, $V_S = 0$

Quartet Resonance in DC Current for $V_a = -V_b$

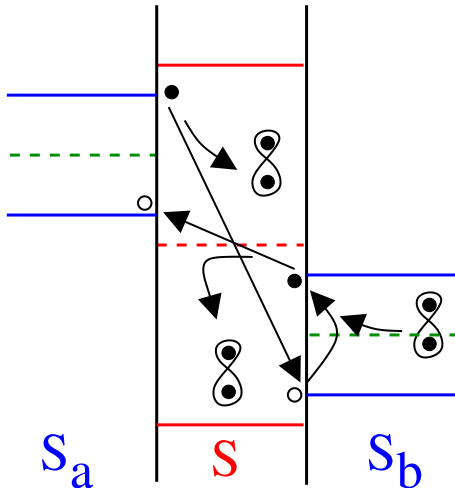
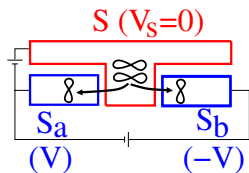
A. Freyn, B. Douçot, D. Feinberg, R. Mélin, PRL 2011, NÉEL / LPTHE



Nonlocal Quartets, DC if $V_a = -V_b$, $V_S = 0$

Quartet Resonance in DC Current for $V_a = -V_b$

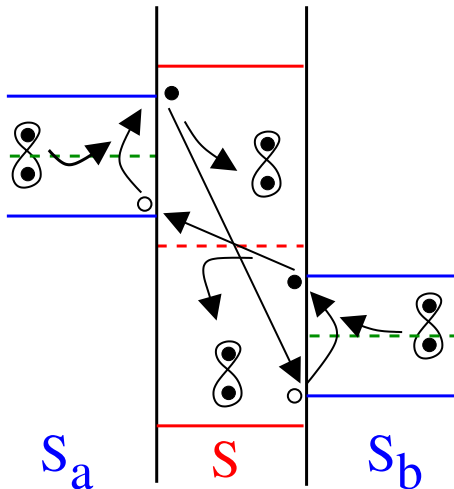
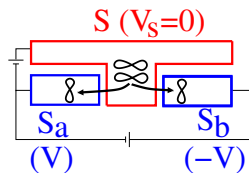
A. Freyn, B. Douçot, D. Feinberg, R. Mélin, PRL 2011, NÉEL / LPTHE



Nonlocal Quartets, DC if $V_a = -V_b$, $V_S = 0$

Quartet Resonance in DC Current for $V_a = -V_b$

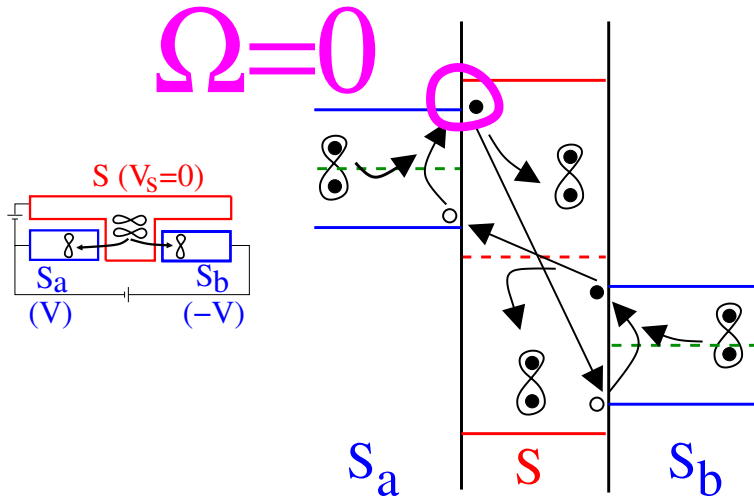
A. Freyn, B. Douçot, D. Feinberg, R. Mélin, PRL 2011, NÉEL / LPTHE



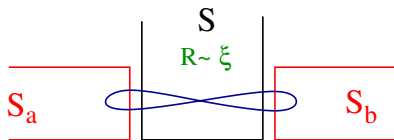
Nonlocal Quartets, DC if $V_a = -V_b$, $V_S = 0$

Quartet Resonance in DC Current for $V_a = -V_b$

A. Freyn, B. Douçot, D. Feinberg, R. Mélin, PRL 2011, NÉEL / LPTHE



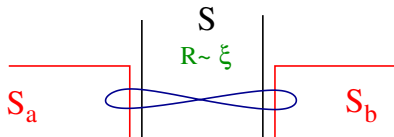
π -junction for the electron quartets



$$\frac{1}{\sqrt{2}} \left(c_{a,\uparrow}^+ c_{b,\downarrow}^+ - c_{a,\downarrow}^+ c_{b,\uparrow}^+ \right)$$

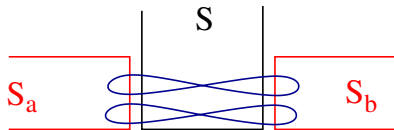
Split pair
Unstable

π -junction for the electron quartets



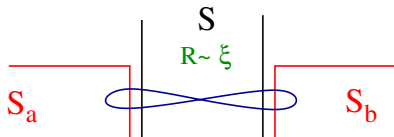
Split pair
Unstable

$$\frac{1}{\sqrt{2}} \left(c_{a,\uparrow}^+ c_{b,\downarrow}^+ - c_{a,\downarrow}^+ c_{b,\uparrow}^+ \right)$$



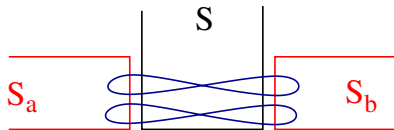
$$\frac{1}{2} \left(c_{a,\uparrow}^+ c_{b,\downarrow}^+ - c_{a,\downarrow}^+ c_{b,\uparrow}^+ \right)^2$$

π -junction for the electron quartets



$$\frac{1}{\sqrt{2}} \left(c_{a,\uparrow}^+ c_{b,\downarrow}^+ - c_{a,\downarrow}^+ c_{b,\uparrow}^+ \right)$$

Split pair
Unstable

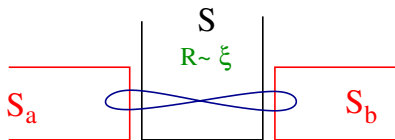


$$\frac{1}{2} \left(c_{a,\uparrow}^+ c_{b,\downarrow}^+ - c_{a,\downarrow}^+ c_{b,\uparrow}^+ \right)^2$$

=

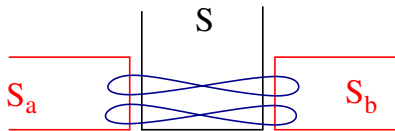
$$- \left(c_{a,\uparrow}^+ c_{a,\downarrow}^+ \right) \left(c_{b,\uparrow}^+ c_{b,\downarrow}^+ \right)$$

π -junction for the electron quartets



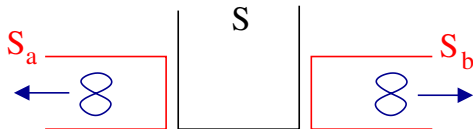
$$\frac{1}{\sqrt{2}} \left(c_{a,\uparrow}^+ c_{b,\downarrow}^+ - c_{a,\downarrow}^+ c_{b,\uparrow}^+ \right)$$

Split pair
Unstable



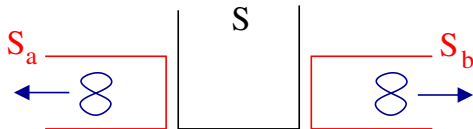
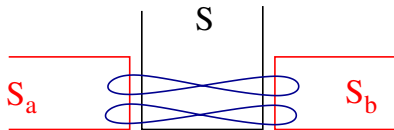
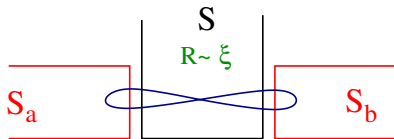
$$\frac{1}{2} \left(c_{a,\uparrow}^+ c_{b,\downarrow}^+ - c_{a,\downarrow}^+ c_{b,\uparrow}^+ \right)^2$$

=



$$- \left(c_{a,\uparrow}^+ c_{a,\downarrow}^+ \right) \left(c_{b,\uparrow}^+ c_{b,\downarrow}^+ \right)$$

π -junction for the electron quartets



$$\frac{1}{2} \left(c_{a,\uparrow}^+ c_{b,\downarrow}^+ - c_{a,\downarrow}^+ c_{b,\uparrow}^+ \right)^2$$

$$= - \left(c_{a,\uparrow}^+ c_{a,\downarrow}^+ \right) \left(c_{b,\uparrow}^+ c_{b,\downarrow}^+ \right)$$

No preformed quartets

Glue between pairs
= interfaces with $R \sim \xi$

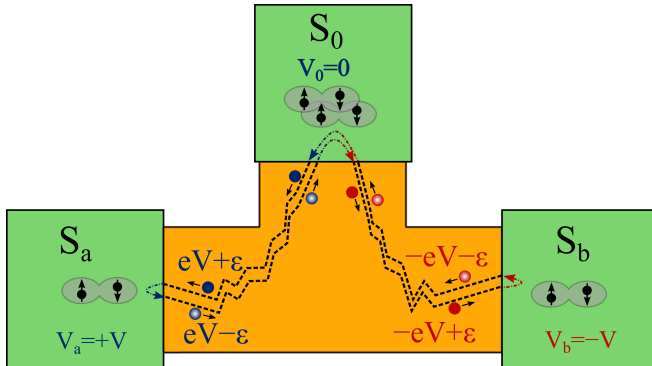
– sign $\Rightarrow \pi$ -junction

$$I_{quartet} = -|I_c| \sin(\varphi_a + \varphi_b - 2\varphi_S)$$

Macroscopic
manifestation of the
internal structure of a
pair (orbital and spin
symmetries)

Experiment on Quartets in a Metallic Cu-Al Structure (Grenoble group)

A.H. Pfeffer, J.E. Duvauchelle, H. Courtois, R. Mélin, D. Feinberg and F. Lefloch,
PRB '14

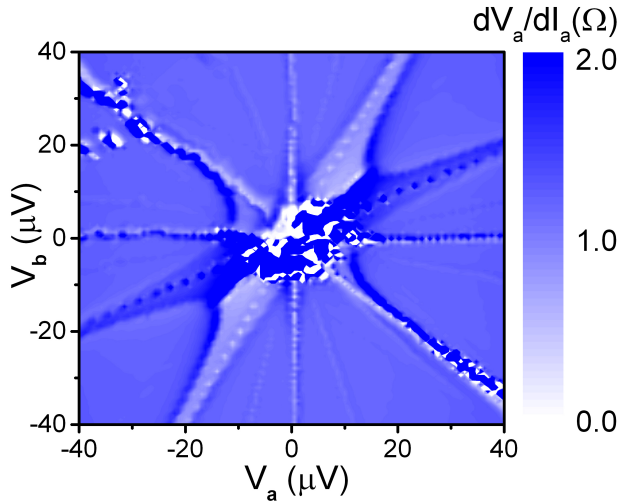
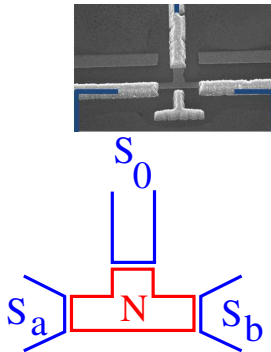


Theoretical calculation

- Perturbative expansion in the tunnel amplitudes
- ⇒ Diffusion modes, evaluated in the ladder approximation

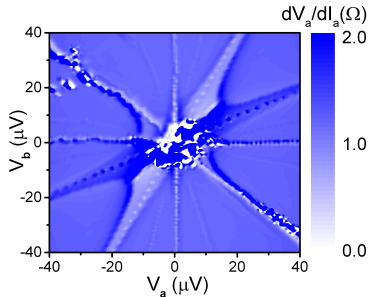
Synchronization for a trijunction ($T = 200$ mK)

A.H. Pfeffer, J.E. Duvauchelle, H. Courtois, F. Lefloch, R. Mélin, D. Feinberg, PRB'14



Synchronization for a trijunction

A.H. Pfeffer, J.E. Duvauchelle, H. Courtois, F. Lefloch, R. Mélin, D. Feinberg, PRB'14

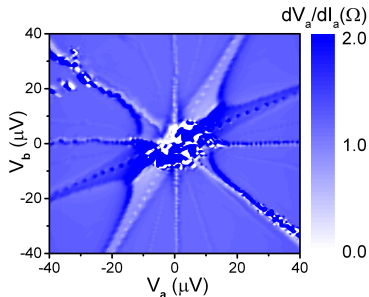


Direct Josephson $S_a - S_0$ and $S_b - S_0$

Two resonances for $V_a = 0$ and $V_b = 0$

Synchronization for a trijunction

A.H. Pfeffer, J.E. Duvauchelle, H. Courtois, F. Lefloch, R. Mélin, D. Feinberg, PRB'14

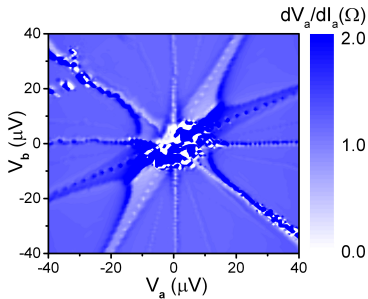


Direct Josephson $S_a - S_b$

- Resonance expected at $V_a = V_b$
- However, not visible because lock-in excitation on S_0

Synchronization for a trijunction

A.H. Pfeffer, J.E. Duvauchelle, H. Courtois, F. Lefloch, R. Mélin, D. Feinberg, PRB'14



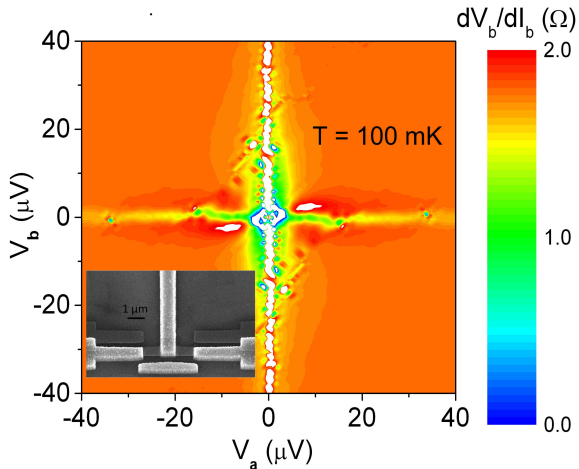
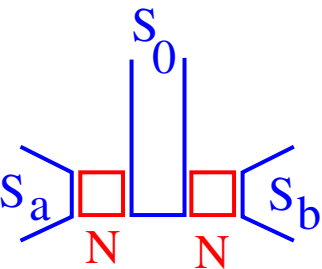
Three additional resonance lines

- $2V_0 = V_a + V_b$; $2V_a = V_0 + V_b$; $2V_b = V_0 + V_a$
- Just permutation of the 3 terminals \rightarrow equivalent resonances
- Are they due to quartets or to classical synchronization by an external impedance ?

No Resonances for a Separated Bijunction

A.H. Pfeffer, J.E. Duvauchelle, H. Courtois, R. Mélin, D. Feinberg, F. Lefloch, PRB '14

2 resonances for direct Josephson $S_0 - S_a$ and $S_0 - S_b$



Zero-Frequency Cross-Correlations

Weizman group experiment

PHYSICAL REVIEW B **93**, 115436 (2016)

Gate-tunable zero-frequency current cross correlations of the quartet state in a voltage-biased three-terminal Josephson junction

Régis Mélin, Moïse Sotto, and Denis Feinberg

*Université Grenoble-Alpes, Institut Néel, BP 166, F-38042 Grenoble Cedex 9, France
and CNRS, Institut Néel, BP 166, F-38042 Grenoble Cedex 9, France*

Jean-Guy Caputo

Laboratoire de Mathématiques, INSA de Rouen, Avenue de l'Université, F-76801 Saint-Etienne du Rouvray, France

Benoît Douçot

*Laboratoire de Physique Théorique et des Hautes Energies, CNRS UMR 7589, Université Pierre et Marie Curie, Sorbonne Universités,
4 Place Jussieu, 75252 Paris Cedex 05*

(Received 26 November 2015; revised manuscript received 3 March 2016; published 25 March 2016)

A three-terminal Josephson junction biased at opposite voltages can sustain a phase-sensitive dc current carrying three-body static phase coherence, known as the “quartet current.” We calculate the zero-frequency current noise cross correlations and answer the question of whether this current is noisy (like a normal current in response to a voltage drop) or noiseless (like an equilibrium supercurrent in response to a phase drop). A quantum dot with a level at energy ϵ_0 is connected to three superconductors S_a , S_b , and S_c with gap Δ , biased at $V_a = V$,

Nonlocal supercurrent of quartets in a three-terminal Josephson junction

Yonatan Cohen^{a,1}, Yuval Ronen^{a,1}, Jung-Hyun Kang^a, Moty Heiblum^{a,2}, Denis Feinberg^b, Régis Mélin^b, and Hadas Shtrikman^a

^aDepartment of Condensed Matter Physics, Braun Center for Submicron Research, Weizmann Institute of Science, 76100 Rehovot, Israel; and ^bInstitut Néel, CNRS, Université Grenoble-Alpes, Institute of Engineering (INP), 38000 Grenoble, France

Edited by Eduardo Fradkin, University of Illinois at Urbana-Champaign, Urbana, IL, and approved May 21, 2018 (received for review January 2, 2018)

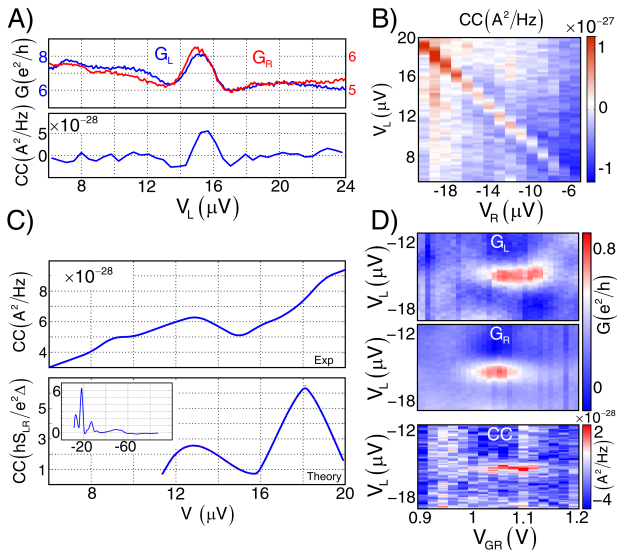
A novel nonlocal supercurrent, carried by quartets, each consisting of four electrons, is expected to appear in a voltage-biased three-terminal Josephson junction. This supercurrent results from a non-

several alternative models for that current could not be ruled out. Here, we verify an emergent coherent quartet supercurrent in a 3TJ, which is formed in a proximitized semiconducting



Zero-Frequency Cross-Correlations

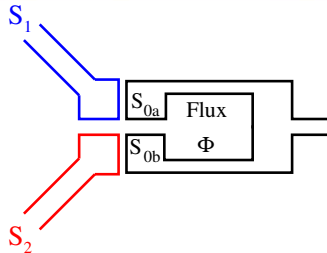
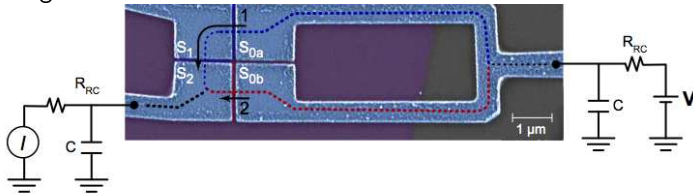
Y. Cohen, Y. Ronen, J.-H. Kang, M. Heiblum,
D. Feinberg, R. Mélin and H. Shtrikman, PNAS 2018



The device

Harvard group experiment

K.-F. Huang, Y. Ronen, R. Mélin, D. Feinberg, K. Watanabe
T. Taniguchi and P. Kim, Nature Comm. '22



The quartet resonance line ($1/2$)

K.-F. Huang, Y. Ronen, R. Mélin, D. Feinberg, K. Watanabe, T. Taniguchi
and P. Kim, Nature Comm. 2022

A. Freyn, B. Douçot, D. Feinberg and R. Mélin,
Phys. Rev. Lett. 106, 257005 (2011)

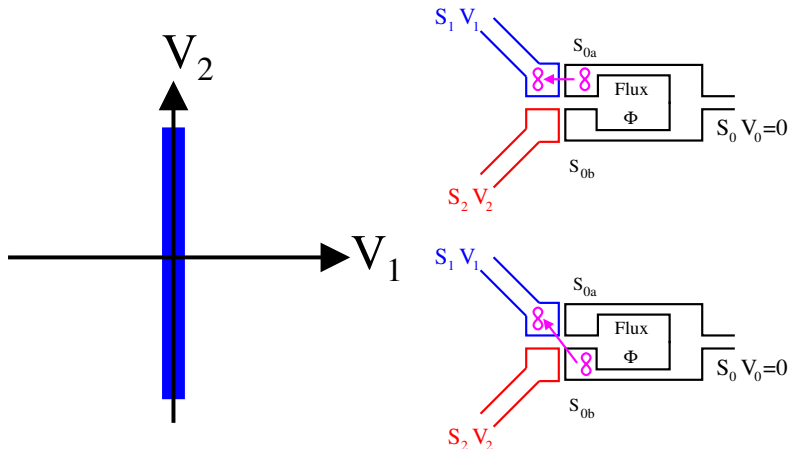
First, the two-terminal processes

The quartet resonance line (1/2)

K.-F. Huang, Y. Ronen, R. Mélin, D. Feinberg, K. Watanabe, T. Taniguchi and P. Kim, Nature Comm. 2022

A. Freyn, B. Douçot, D. Feinberg and R. Mélin, Phys. Rev. Lett. 106, 257005 (2011)

DC-Josephson effect between S_1 and $S_0 \Rightarrow V_1 = 0$

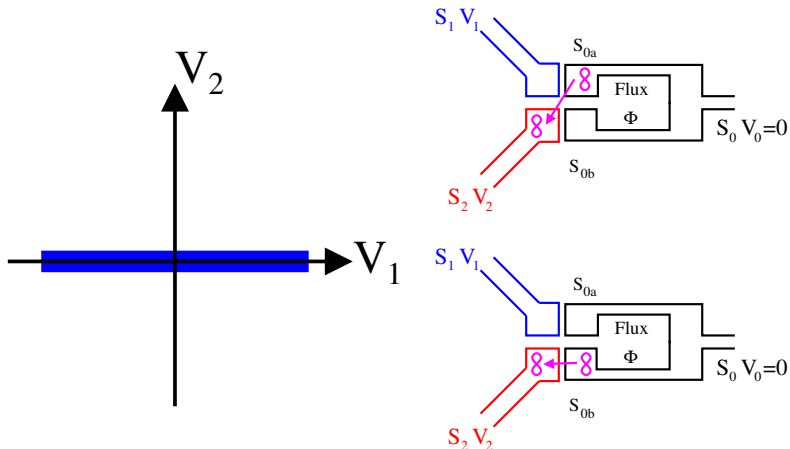


The quartet resonance line (1/2)

K.-F. Huang, Y. Ronen, R. Mélin, D. Feinberg, K. Watanabe, T. Taniguchi and P. Kim, Nature Comm. 2022

A. Freyn, B. Douçot, D. Feinberg and R. Mélin, Phys. Rev. Lett. 106, 257005 (2011)

DC-Josephson effect between S_2 and $S_0 \Rightarrow V_2 = 0$

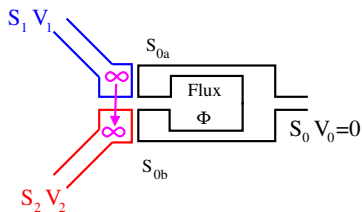
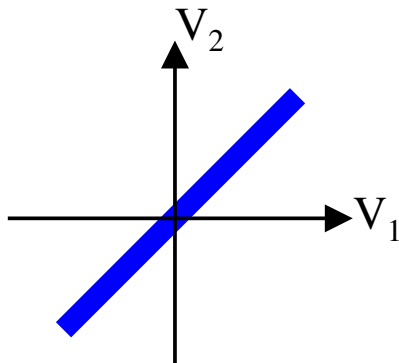


The quartet resonance line (1/2)

K.-F. Huang, Y. Ronen, R. Mélin, D. Feinberg, K. Watanabe, T. Taniguchi and P. Kim, Nature Comm. 2022

A. Freyn, B. Douçot, D. Feinberg and R. Mélin, Phys. Rev. Lett. 106, 257005 (2011)

DC-Josephson effect between S_1 and $S_2 \Rightarrow V_1 = V_2$



The quartet resonance line (1/2)

K.-F. Huang, Y. Ronen, R. Mélin, D. Feinberg, K. Watanabe, T. Taniguchi
and P. Kim, Nature Comm. 2022

A. Freyn, B. Douçot, D. Feinberg and R. Mélin,
Phys. Rev. Lett. 106, 257005 (2011)

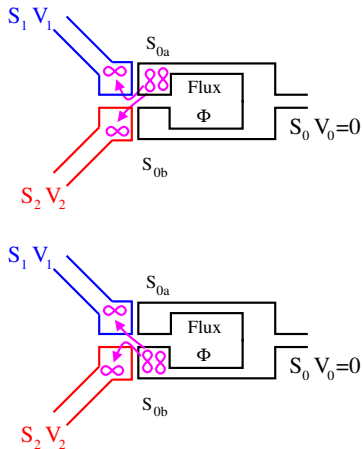
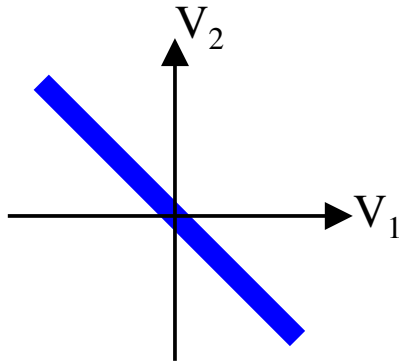
Second, the three-terminal processes

The quartet resonance line (1/2)

K.-F. Huang, Y. Ronen, R. Mélin, D. Feinberg, K. Watanabe, T. Taniguchi and P. Kim, Nature Comm. 2022

A. Freyn, B. Douçot, D. Feinberg and R. Mélin, Phys. Rev. Lett. 106, 257005 (2011)

Quartet from $S_{0,a}$ or $S_{0,b}$ towards $(S_1, S_2) \Rightarrow V_1 + V_2 = 0$



The quartet resonance line ($1/2$)

K.-F. Huang, Y. Ronen, R. Mélin, D. Feinberg, K. Watanabe, T. Taniguchi
and P. Kim, Nature Comm. 2022

A. Freyn, B. Douçot, D. Feinberg and R. Mélin,
Phys. Rev. Lett. 106, 257005 (2011)

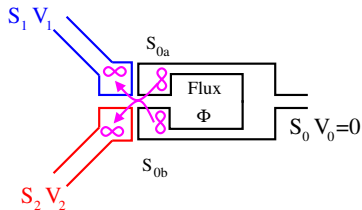
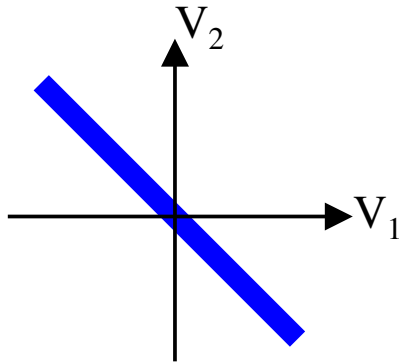
Third, the four-terminal processes

The quartet resonance line (1/2)

K.-F. Huang, Y. Ronen, R. Mélin, D. Feinberg, K. Watanabe, T. Taniguchi and P. Kim, Nature Comm. 2022

A. Freyn, B. Douçot, D. Feinberg and R. Mélin, Phys. Rev. Lett. 106, 257005 (2011)

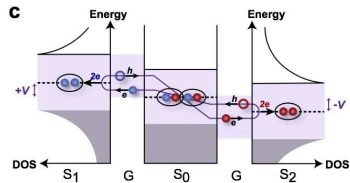
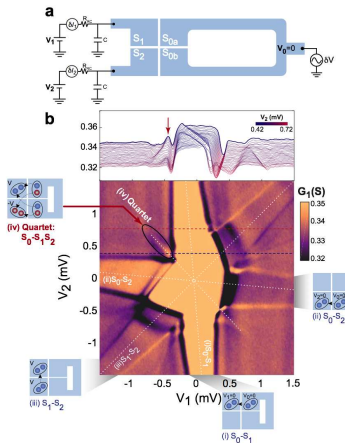
Split quartet from $(S_{0,a}, S_{0,b})$ towards $(S_1, S_2) \Rightarrow V_1 + V_2 = 0$



The quartet resonance line (2/2)

K.-F. Huang, Y. Ronen, R. Mélin, D. Feinberg, K. Watanabe, T. Taniguchi and P. Kim, Nature Comm. 2022

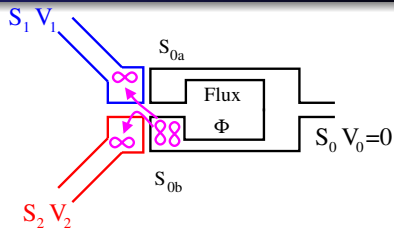
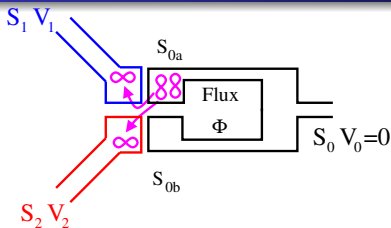
A. Freyn, B. Douçot, D. Feinberg and R. Mélin, Phys. Rev. Lett. 106, 257005 (2011)



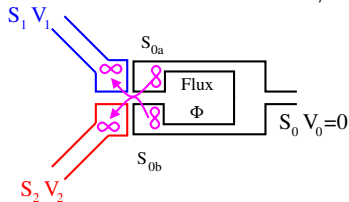
The flux periodicity (1/2)

K.-F. Huang, Y. Ronen, R. Mélin, D. Feinberg, K. Watanabe, T. Taniguchi and P. Kim, Nature Comm. 2022

R. Mélin, Phys. Rev. B **102**, 245435 (2020)



Interference between $3TQ_a$, $3TQ_b \Rightarrow h/4e$ -periodicity



Interference between $3TQ_a$, $3TQ_b$, $4TSQ \Rightarrow h/2e$ -periodicity

The flux periodicity (2/2)

K.-F. Huang, Y. Ronen, R. Mélin, D. Feinberg, K. Watanabe, T. Taniguchi and P. Kim, Nature Comm. 2022

In experiments, balance between the $3TQ$ and the $4TSQ$ tuned by gate voltage.

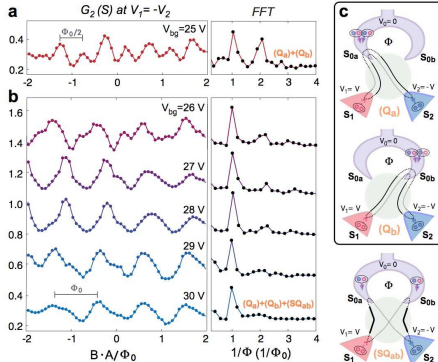
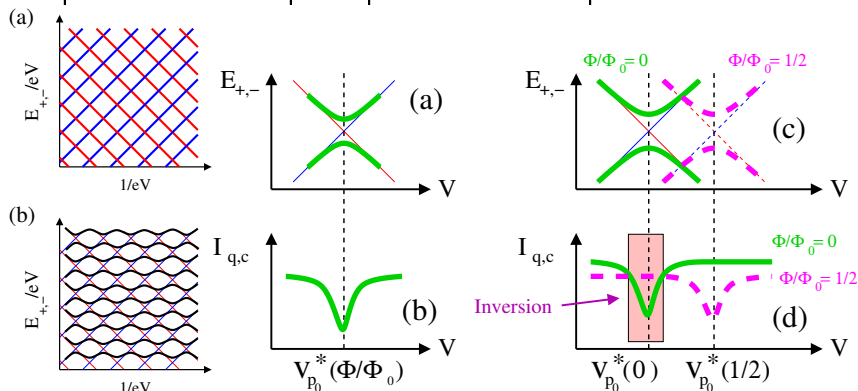


Figure 3. Different types of quartet process. **a**, Left panel shows the quartet differential conductance $G_2 (= \delta I_2 / \delta V)$ measured at S_2 as a function of the flux $\Phi = B \cdot A / \Phi_0$ for $V_{bg} = 25$ V. Satisfying the quartet bias condition, $V_1 = -V_2$ is fixed at 4 V. Right panel shows the analysis from fast Fourier transform (FFT) with a prominent second harmonic. **b**, for $V_{bg} = 26$ –30 V. The periodicity evolves from half-flux quantum to one flux quantum as V_{bg} increases. **c**, Subfigure (Q_a) (Q_b) shows the conventional three-terminal quartet process with only one out of the two loop contacts involved. Electron-hole conversion happens twice at the same contact of the loop (either S_{0a} or S_{0b}), resulting in periodicity of half-flux quantum. Subfigure (SQ_{ab}) shows the split-quartet process involving both contacts of the loop. With the odd parity of Cooper pairs transferred, the periodicity is one flux quantum.

Floquet resonances and the voltage dependence

R. Mélin and B. Douçot, Phys. Rev. B **102**, 245436 (2020)

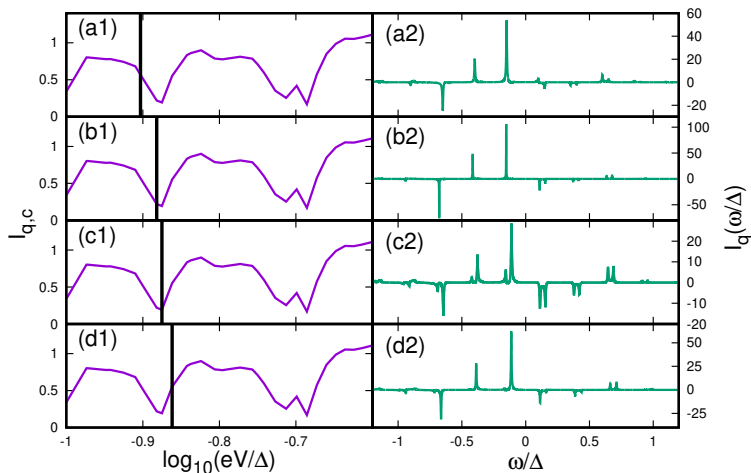
0D quantum dots. Floquet replica of the ABS spectrum.



Hybridization between the two Floquet ladders

R. Mélin and B. Douçot, Phys. Rev. B **102**, 245436 (2020)

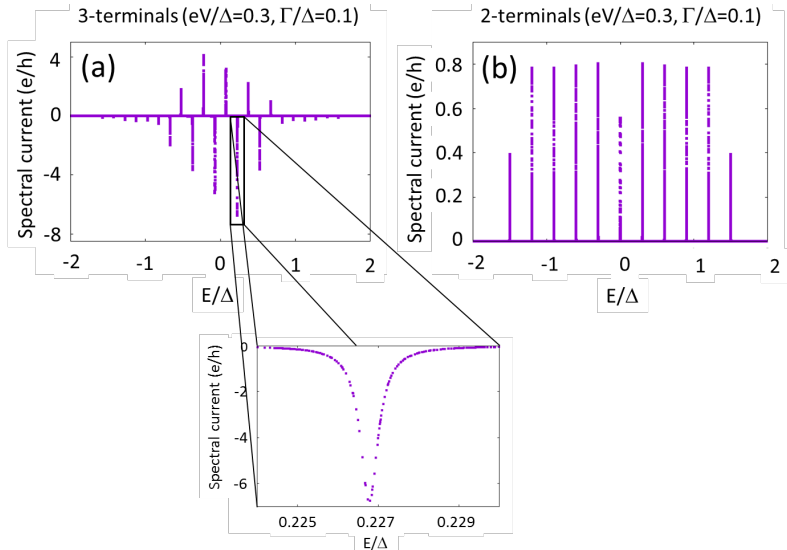
$$\Phi/\Phi_0=0, \gamma/\Delta=0.3$$



Spectral Current:

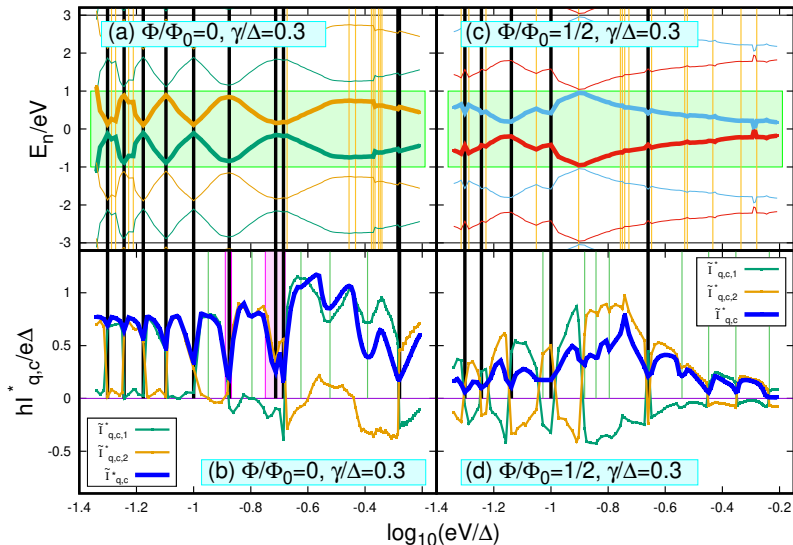
Evidence for two Floquet-Wannier-Stark Ladders

R. Mélin, J.-G. Caputo, K. Yang and B. Douçot, PRB '17



Quartet critical current-voltage characteristics

R. Mélin and B. Douçot, Phys. Rev. B **102**, 245436 (2020)



Previous experiments were shown to be compatible with the quartets:

- The Grenoble group experiment:

A. H. Pfeffer, J. E. Duvauchelle, H. Courtois, R. Mélin, D. Feinberg, and F. Lefloch,
Phys. Rev. B **90**, 075401 (2014).

- The Weizmann group experiment:

Y. Cohen, Y. Ronen, J.-H. Kang, M. Heiblum, D. Feinberg, R. Mélin, and H. Shtrikman,
PNAS July **3**, 2018 115 (27) 6991-6994.

Compatibility between new Harvard group experiment and theory of four-terminal quartets.

Conclusions (2/2)

- Interference between quartet-ABS produces cross-over between $h/4e$ - and $h/2e$ -periodicity.
- New channel of four-terminal split quartets produces $I_{q,c}(\Phi = \pi, V) \neq I_{q,c}(\Phi = 0, V)$.

Two mechanisms for Floquet theory:

- Without relaxation: dips in the quartet critical current-voltage characteristics due to quantum coherent superpositions of the time-dependent ABS.
- With relaxation: Resonance peaks are inverse proportional to the Floquet line-width.